

# 14 A Pseudospectral method for periodic functions

**Lab Objective:** We look at a pseudospectral method with a Fourier basis, and numerically solve the advection equation using a pseudospectral discretization in space and a Runge-Kutta integration scheme in time.

Let  $f$  be a periodic function on  $[0, 2\pi]$ . Let  $x_1, \dots, x_N$  be  $N$  evenly spaced grid points on  $[0, 2\pi]$ . Since  $f$  is periodic on  $[0, 2\pi]$ , we can ignore the grid point  $x_N = 2\pi$ . We will further assume that  $N$  is even; similar formulas can be derived for  $N$  odd. Let  $h = 2\pi/N$ ; then  $\{x_0, \dots, x_{N-1}\} = \{0, h, 2h, \dots, 2\pi - h\}$ .

The discrete Fourier transform (DFT) of  $f$ , denoted by  $\hat{f}$  or  $\mathcal{F}(f)$ , is given by

$$\hat{f}(k) = h \sum_{j=0}^{N-1} e^{-ikx_j} f(x_j) \quad \text{where } k = -N/2 + 1, \dots, 0, 1, \dots, N/2.$$

The inverse DFT is then given by

$$f(x_j) = \frac{1}{2\pi} \sum_{k=-N/2}^{N/2} \frac{e^{ikx_j}}{c_k} \hat{f}(k), \quad j = 0, \dots, N-1, \quad (14.1)$$

where

$$c_k = \begin{cases} 2 & \text{if } k = -N/2 \text{ or } k = N/2, \\ 1 & \text{otherwise.} \end{cases} \quad (14.2)$$

The inverse DFT can then be used to define a natural interpolant (sometimes called a band-limited interpolant) by evaluating (14.1) at any  $x$  rather than  $x_j$ :

$$p(x) = \frac{1}{2\pi} \sum_{k=-N/2}^{N/2} e^{ikx} \hat{f}(k). \quad (14.3)$$

The interpolant for  $f'$  is then given by

$$p'(x) = \frac{1}{2\pi} \sum_{k=-N/2+1}^{N/2-1} ik e^{ikx} \hat{f}(k). \quad (14.4)$$

Consider the function  $u(x) = \sin^2(x) \cos(x) + e^{2\sin(x+1)}$ . Using (14.4), the derivative  $u'$  may be approximated with the following code.<sup>1</sup> We note that although we only approximate  $u'$  at the Fourier grid points, (14.4) provides an analytic approximation of  $u'$  in the form of a trigonometric polynomial.

```
import numpy as np
from scipy.fftpack import fft, ifft
import matplotlib.pyplot as plt

N=24
x1 = (2.*np.pi/N)*np.arange(1,N+1)
f = np.sin(x1)**2.*np.cos(x1) + np.exp(2.*np.sin(x1+1))

# This array is reordered in Python to
# accomodate the ordering inside the fft function in scipy.
k = np.concatenate(( np.arange(0,N/2) ,
                     np.array([0]) , # Because hat{f}'(k) at k = N/2 is zero.
                     np.arange(-N/2+1,0,1) ))

# Approximates the derivative using the pseudospectral method
f_hat = fft(f)
fp_hat = ((1j*k)*f_hat)
fp = np.real(ifft(fp_hat))

# Calculates the derivative analytically
x2 = np.linspace(0,2*np.pi,200)
derivative = (2.*np.sin(x2)*np.cos(x2)**2. -
             np.sin(x2)**3. +
             2*np.cos(x2+1)*np.exp(2*np.sin(x2+1))
             )

plt.plot(x2,derivative, '-k',linewidth=2.)
plt.plot(x1,fp, '*b')
plt.savefig('spectral2_derivative.pdf')
plt.show()
```

**Problem 1.** Consider again the function  $u(x) = \sin^2(x) \cos(x) + e^{2\sin(x+1)}$ . Create a function that approximates  $\frac{1}{2}u'' - u'$  on the Fourier grid points for  $N = 24$ .

## The advection equation

Recall that the advection equation is given by

$$u_t + cu_x = 0 \tag{14.5}$$

<sup>1</sup>See *Spectral Methods in MATLAB* by Lloyd N. Trefethen. Another good reference is *Chebyshev and Fourier Spectral Methods* by John P. Boyd.

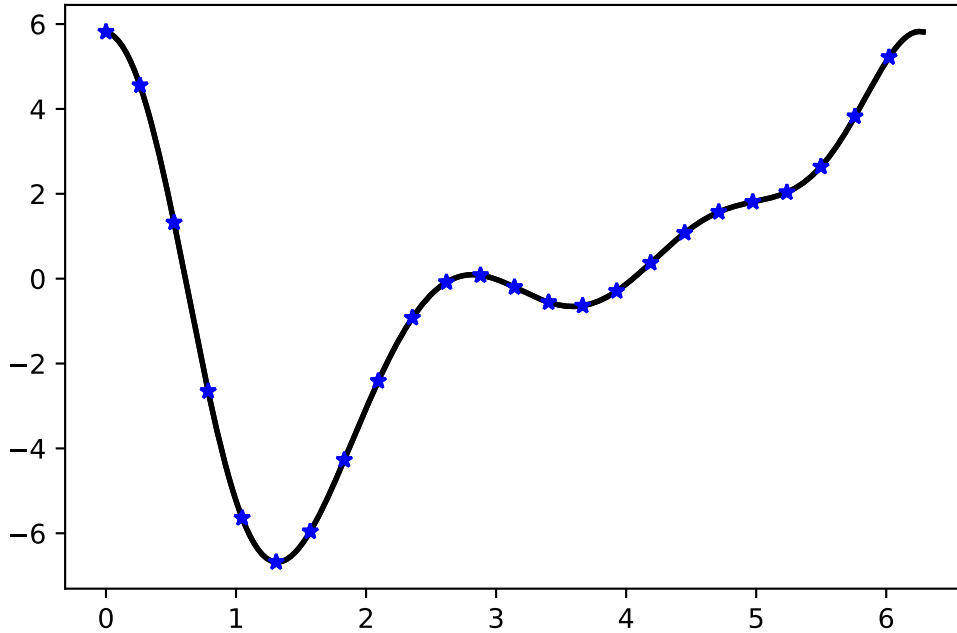


Figure 14.1: The derivative of  $u(x) = \sin^2(x) \cos(x) + e^{2 \sin(x+1)}$ .

where  $c$  is the speed of the wave (the wave travels to the right for  $c > 0$ ). We will consider the solution of the advection equation on the circle; this essentially amounts to solving the advection equation on  $[0, 2\pi]$  and assuming periodic boundary conditions.

A common method for solving time-dependent PDEs is called the *method of lines*. To apply the method of lines to our problem, we use our Fourier grid points in  $[0, \pi]$ : given an even  $N$ , let  $h = 2\pi/N$ , so that  $\{x_0, \dots, x_{N-1}\} = \{0, h, 2h, \dots, 2\pi - h\}$ . By using these grid points we obtain the collection of equations

$$u_t(x_j, t) + cu_x(x_j, t) = 0, \quad t > 0, \quad j = 0, \dots, N-1. \quad (14.6)$$

Let  $U(t)$  be the vector valued function given by  $U(t) = (u(x_j, t))_{j=0}^{N-1}$ . Let  $\mathcal{F}(U)(t)$  denote the discrete Fourier transform of  $u(x, t)$  (in space), so that

$$\mathcal{F}(U)(t) = (\hat{u}(k, t))_{k=-N/2+1}^{N/2}.$$

Define  $\mathcal{F}^{-1}$  similarly. Using the pseudospectral approximation in space leads to the system of ODEs

$$U_t + \vec{c}\mathcal{F}^{-1} \left( i\vec{k}\mathcal{F}(U) \right) = 0 \quad (14.7)$$

where  $\vec{k}$  is a vector, and  $\vec{k}\mathcal{F}(U)$  denotes element-wise multiplication. Similarly  $\vec{c}$  could also be a vector, if the wave speed  $c$  is allowed to vary.

**Problem 2.** Using a fourth order Runge-Kutta method (RK4), solve the initial value problem

$$u_t + c(x)u_x = 0, \quad (14.8)$$

where  $c(x) = .2 + \sin^2(x - 1)$ , and  $u(x, t = 0) = e^{-100(x-1)^2}$ . Plot your numerical solution from  $t = 0$  to  $t = 8$  over 150 time steps and 100  $x$  steps. Note that the initial data is nearly zero near  $x = 0$  and  $2\pi$ , and so we can use the pseudospectral method. <sup>a</sup> Use the following code to help graph.

```
t_steps = 150    # Time steps
x_steps = 100    # x steps

'''
Your code here to set things up
'''

sol = # RK4 method. Should return a t_steps by x_steps array

X,Y = np.meshgrid(x_domain, t_domain)
fig = plt.figure()
ax = fig.add_subplot(111, projection="3d")
ax.plot_wireframe(X,Y,sol)
ax.set_zlim(0,3)
plt.show()
```

<sup>a</sup>This problem is solved in *Spectral Methods in MATLAB* using a leapfrog discretization in time.

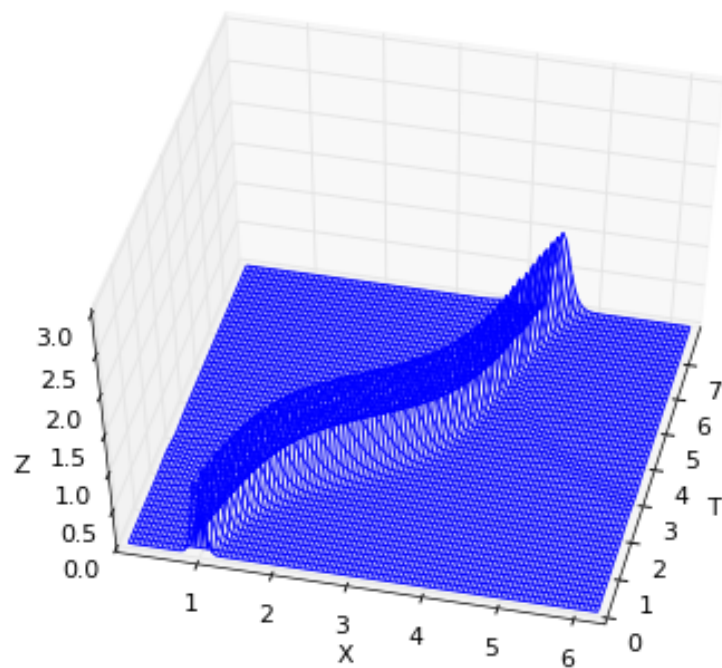


Figure 14.2: The solution of the variable speed advection equation; see Problem 2.